Physical Principles in Biology Biology 3550 Spring 2024

Lecture 24

Thermodynamics:

The First Law, State Functions and Free Energy

Wednesday, 13 March 2024

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Announcements

Midterm Exam:

- Friday, 15 March
- Will cover all material from before Spring Break
- 50 min
- Review Session
 - 5:15 PM, Thursday, 14 March
 - HEB 2002
 - Come with questions!

Adiabatic (without heat flow) Expansion of a Gas



Adiabatic Gas Expansion With Work



Isothermal Expansion with Work



Keeping Score

Change in energy of the gas molecules (the "system"):

$$\Delta E = E_{\text{final}} - E_{\text{start}}$$

Work, w:

- w > 0, when work is done <u>on</u> the system.
- w < 0, when the system does work on the outside world, as in the expansion of the gas.
- Heat, q.
 - *q* > 0, when heat flows into the system.
 - q < 0, when heat flows out of the system into the surroundings.
- For both work, *w*, and heat, *q*, a positive value indicates a transfer to the system from the surroundings.

The First Law of Thermodynamics

- Common statements in words:
 - "The energy of the universe is conserved"
 - "Energy cannot be created or destroyed" s
 - Later modified to account for interconversion of mass and energy. (Einstein's $E = mc^2$)
- The formal mathematical statement: For any process,

 $\Delta E = q + w$

- Any change in the energy of the system has to be accounted for by work or heat.
- Work and heat represent the transfer of energy from the surroundings to the system.
- Ignores other forms of energy, such as electromagnetic radiation.



A)
$$q < 0$$

C) q > 0



■ *q* < 0

■ *q* > 0

a = 0

Adiabatic Expansion Without Work

Perfect insulation

Which of the following is true?

B)
$$w = 0$$

C) $w > 0$









A) q < 0

B) q = 0

C)



Which of the following is true?





State Functions and Path-dependent Functions

State functions of a system depend only on the current state of the system and do not depend on history.

Examples:

- Temperature, T
- Pressure, P
- Volume, V
- Energy, E
- For any change in a system, the change in a state function depends only on the beginning and ending states.

$$\Delta T = T_{\text{final}} - T_{\text{start}}$$

 $\Delta P = P_{\text{final}} - P_{\text{start}}$
 $\Delta V = V_{\text{final}} - V_{\text{start}}$
 $\Delta E = E_{\text{final}} - E_{\text{start}}$

Work, *w*, and heat, *q*, are <u>not</u> state functions.

Adiabatic vs. Isothermal Gas Expansion



The two processes start and end in the same states:

$$\Delta V_{ad} = \Delta V_{isot} > 0$$
 $\Delta P_{ad} = \Delta P_{isot} < 0$
 $\Delta T_{ad} = \Delta T_{isot} = 0$ $\Delta E_{ad} = \Delta E_{isot} = 0$

But, the heat and work for the two processes are different:

$$w_{
m ad} = 0$$
 $w_{
m isot} < 0$ $q_{
m ad} = 0$ $q_{
m isot} > 0$

Heat and work are "path-dependent" functions.

The Maximum-work Path for Gas Expansion



- Piston is allowed to move upward in infinitesimally small steps.
- Temperature is never allowed to drop.
- Pressure drops as gas expands, so less work is done per step.
- If larger steps are ever taken:
 - The temperature drops.
 - The pressure drops more than it would in an infinitesimal step.
 - Less work is produced.

The Minimum-work Path for Gas Compression



- Piston is pushed down in infinitesimally small steps.
- Energy is transferred from piston to gas molecules.
- Temperature is never allowed to increase.
- Excess energy flows to the reservoir as heat.
- P increases as gas is compressed, so more work is required per step.
- If larger steps are ever taken, more work is required.

A Reversible Cycle of Compression and Expansion



- Steps in both directions are infinitesimal.
- Compression and expansion are exactly the reverse of one another.

$$w_{ ext{comp}} = -w_{ ext{exp}} = -q_{ ext{comp}} = q_{ ext{exp}}$$

For the complete cycle:

$$\Delta E=0 \quad w=0 \quad q=0$$

Either compression or expansion can be reversed at any point by an infinitesimal force in the opposite direction.

- w and q are <u>not</u> state functions.
- BUT, we can define the value of w (or q) for a specific process linking two states to be a change in a state function.
- We define the work for the reversible (infinitely slow) conversion of one state to the another, w_{rev} , to be the change in the state function *F*.
 - ΔF is called the change in "free energy." (Helmholtz free energy)
 - ΔF is the maximum amount of work that can be obtained from the change in state.
 - ΔF is "free energy" in the sense that it is energy that is available to do work.
 - "Free energy" is really the most expensive kind of energy! It's the kind that we pay for to do work.
- The heat for a reversible process, q_{rev}, is also a state function that we will return to shortly.

Calculating the Work for the Reversible Isothermal Expansion of a Gas



The work is calculated by integrating force with respect to distance.

$$w_{\rm rev} = \int_{x_1}^{x_2} f dx$$

f is force, and x represents the position of the piston.

The force can be expressed in terms of pressure, P, and the area that the piston presents to the gas, A:

$$f = -P \cdot A$$

Calculating the Work for the Reversible Isothermal Expansion of a Gas



From the previous slide:

$$w_{\mathsf{rev}} = \int_{x_1}^{x_2} f dx = -\int_{x_1}^{x_2} P \cdot A dx$$

For each small increment in x, there is a corresponding increment in volume, dV.

$$dV = Adx$$

Calculating the Work for the Reversible Isothermal Expansion of a Gas



Substituting dV for Adx in the work integral:

$$w_{\rm rev} = -\int_{x_1}^{x_2} P \cdot A dx = -\int_{V_1}^{V_2} P dV$$

(The product *PV* has the units of energy or work!)

From the ideal gas law, P = nRT/V. Substituting:

$$w_{\rm rev} = -\int_{V_1}^{V_2} \frac{nRT}{V} dV = -nRT \int_{V_1}^{V_2} \frac{dV}{V}$$

Calculating the Work for the Reversible Isothermal Expansion of a Gas



From the previous slide:

$$w_{\rm rev} = -nRT \int_{V_1}^{V_2} \frac{dV}{V}$$

From calculus:

$$w_{\text{rev}} = -nRT \ln(V)\Big|_{V_1}^{V_2} = -nRT \ln\left(\frac{V_2}{V_1}\right)$$

• w_{rev} depends only on n, T, V_1 and V_2 . w_{rev} represents a change in a state function, ΔF .

Calculate the work, w_{rev} , for the reversible expansion of 0.5 mole of gas from 1 L to 5 L at 37° C. $R = 0.082 \,\mathrm{L} \cdot \mathrm{atm} \cdot \mathrm{K}^{-1} \mathrm{mol}^{-1} = 8.3 \,\mathrm{J} \cdot \mathrm{K}^{-1} \mathrm{mol}^{-1}$ A) -2100 J **B**) -900 J C) -250 J **D**) 250 J E) 900 J

$$w_{
m rev} = -nRT \ln \left(rac{V_2}{V_1}
ight) = -0.5 \, {
m mol} imes 8.3 \, {
m J} \cdot {
m K}^{-1} {
m mol}^{-1} imes 310 \, {
m K} imes {
m ln} \left(rac{5 \, {
m L}}{1 \, {
m L}}
ight)$$

Calculating the Work for the Reversible Isothermal Expansion of a Gas



Since the number of moles remains constant, the change in volume can also be expressed in terms of a change in concentration:

$$C_1V_1$$
 = no. of moles = C_2V_2

 $V_2/V_1 = C_1/C_2$

Substituting into the result from the previous slide:

$$w_{\rm rev} = -nRT \ln\left(\frac{C_1}{C_2}\right)$$

$q_{\rm rev}$ for Isothermal Expansion

For an isothermal process (regardless of path):

$$\Delta E = q + w = 0$$

q = -w

For reversible isothermal expansion:

$$q_{
m rev} = -w_{
m rev} = nRT \ln \left(rac{V_2}{V_1}
ight)$$

Sign check:

For expansion, $V_2 > V_1$ and $q_{rev} > 0$, and the heat flow is <u>into</u> the system, as we expect.

For any two states, there is, in principle, a reversible (infinitely slow) process separating them, and we can define q_{rev} and w_{rev}.

Reconsider the Adiabatic Expansion Without Work



- The energy available for work has decreased.
- We suspect that this has to do with the loss of order, or increase in entropy.
- But, what is entropy? How do we give it a number?